



# Tomographic reconstruction for multi-guidestar adaptive optics

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# Fundamental questions for the architecture of AO for extremely large telescopes



- Number of DMs, and their conjugate locations

Tokovinin & LeLouarn, "Isoplanatism in MCAO," JOSA A, Oct 2000

- Number of actuators per DM

Pareto-optimal solution (Gavel, DeKany, Baumann, Nelson)

- Number and placement of *laser* guide stars

## -This talk-

- Brightness of guide stars

Tokovinin & Viard, "Limiting precision...", JOSA A, Apr 2001

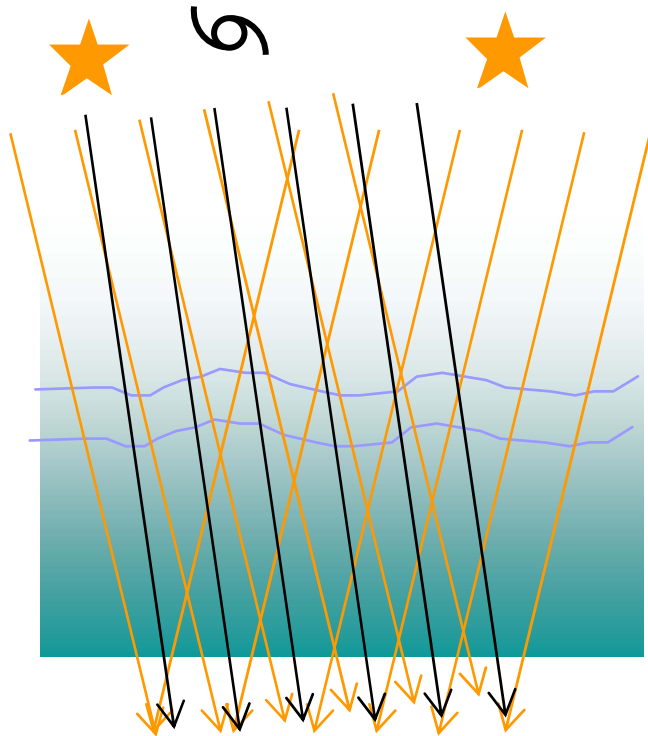
- Controller bandwidth



# Tokovinin's minimum variance wavefront reconstruction algorithm



Tokovinin & Viard, JOSA-A, 18, 4, 2001



1. Measure wavefronts from guidestars at field angles  $\theta_i$ ,  $i=1 \dots \#gs$

$$\phi_i(x) = \int_0^{\infty} \Delta n(x - \theta_i z, z) dz$$

2. Estimate wavefront as it would appear coming from direction  $\theta$  using a minimum variance estimator in the spatial frequency domain

$$\Phi(k_x) = \sum_{\#gs} g(k_x, \theta_i) \Phi_i(k_x)$$

3. Apply wavefront correction for science object at field angle  $\theta$



# Error analysis of Tokovinin's algorithm



**Wavefront error**  
spatial frequency  $k_x$ ,  
science direction  $\theta$

=

Filtered integral of  
index variation,  $\Delta n$ ,  
along path

+

Noise on each WFS,  
filtered by the estimator,  $g_k$

$$\varepsilon(k_x, \theta) = \int P(k_x, z, \theta) \Delta n(k_x, z) dz - \sum_{l=0}^{\#gs} g_l(k_x, \theta) v_l(k_x)$$

$$P(k_x, \theta, z) = \left[ e^{-ik_x \theta z} - \sum_{l=1}^{\#gs} g_l(k_x, \theta) M(k_x) e^{-ik_x \theta_l z} \right]$$

Wavefront sensor  
operator,  $M$

- Choose feedback gains  $g_l(k_x, \theta)$  to minimize the variance of wavefront error, for given  $C_n^2(z) = \langle \Delta n^2(z) \rangle$  and  $\sigma_n^2 = \langle v_l^2 \rangle$

The algorithm  $g_l(k_x, \theta)$  assumes

- infinite aperture
- plane waves



# Tomographic reconstruction error (assumes infinite aperture and plane waves)



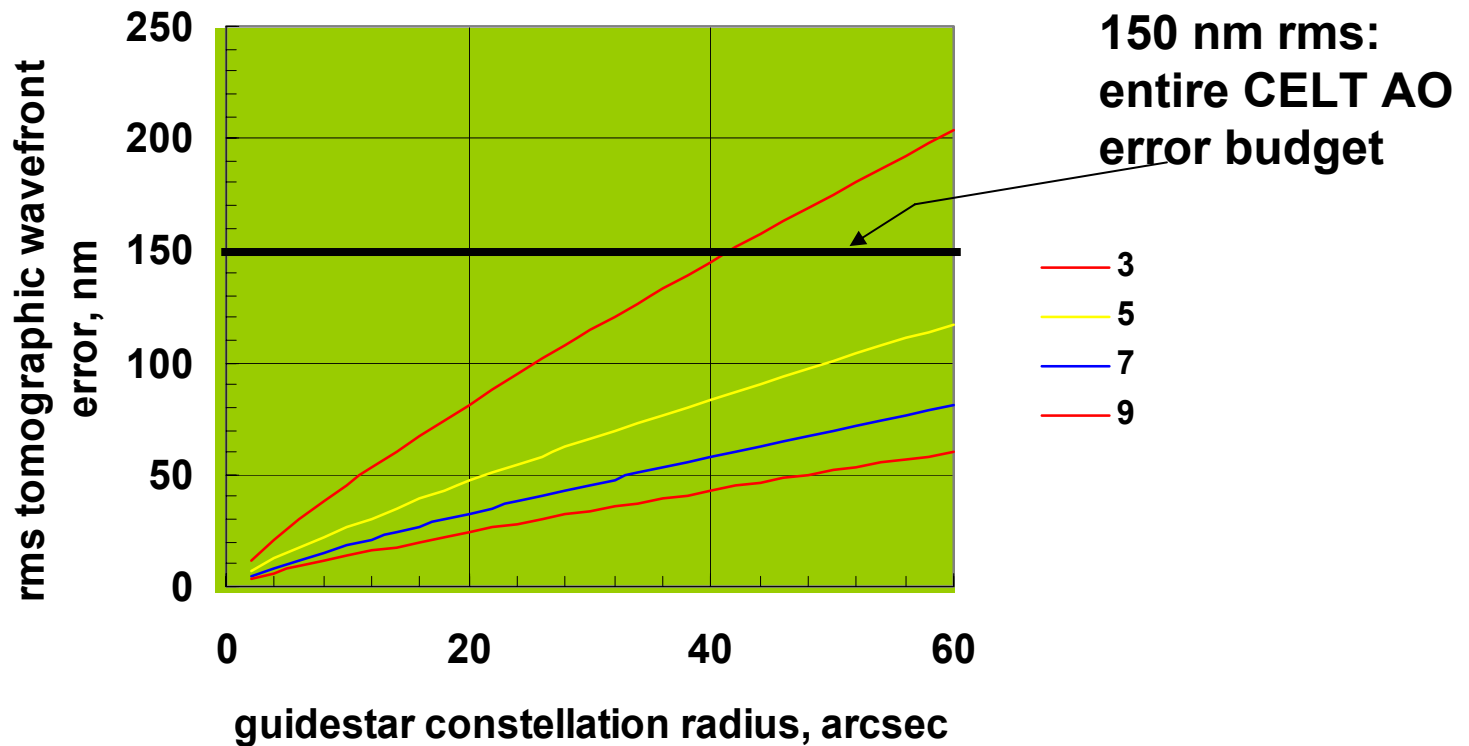
$$\sigma^2 \approx \left( \frac{\Theta \delta}{r_0} \right)^{5/3} e(\theta)$$

$\Theta$  = constellation radius

$r_0$  = transverse coherence distance (Fried's parameter)

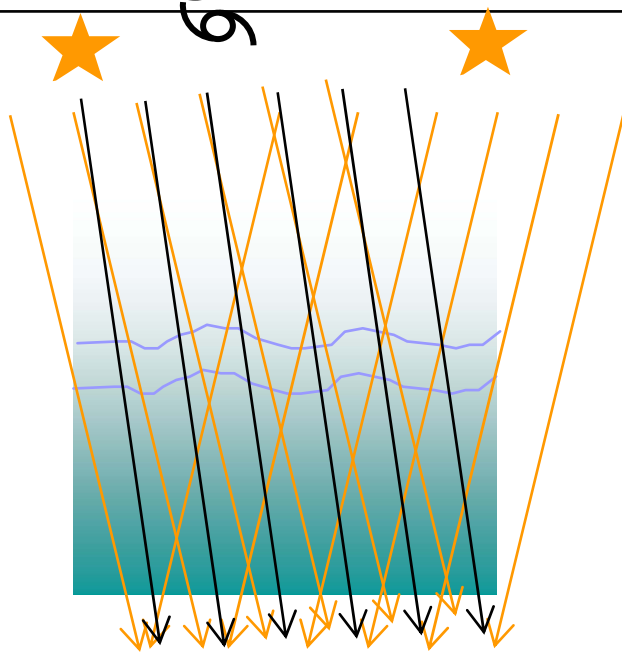
$\delta$  = effective layer thickness

$e$  = field-dependent factor ( $\leq 1$  inside constellation)

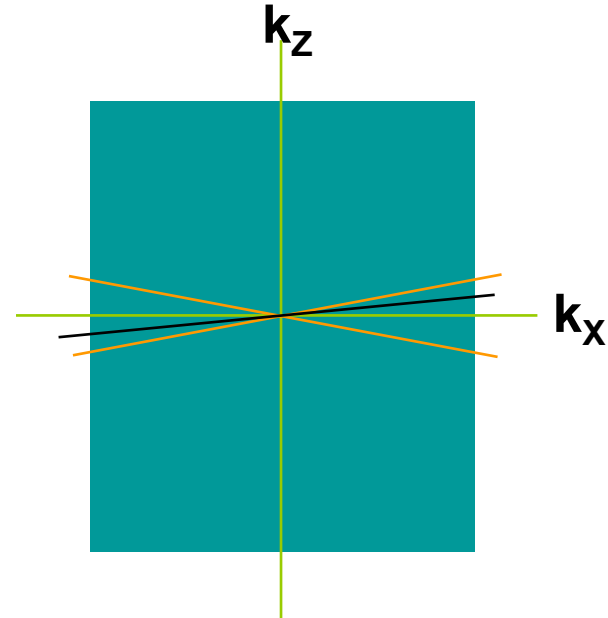




# Fourier transform interpretation of tomographic wavefront reconstruction



$$\phi(x) = \int_0^{\infty} \Delta n(x - \theta z, z) dz$$



$$\Phi(k_x) = \Delta N(k_x, -k_x \theta)$$

## Fourier slice theorem in tomography (Kak, 1988)

- Each wavefront sensor measures the integral of index variation along the ray lines
- The line integral along z determines the  $k_z=0$  Fourier spatial frequency component
- Projections at several angles sample the  $k_x, k_y, k_z$  volume



# Consequences of finite number and field of guide stars

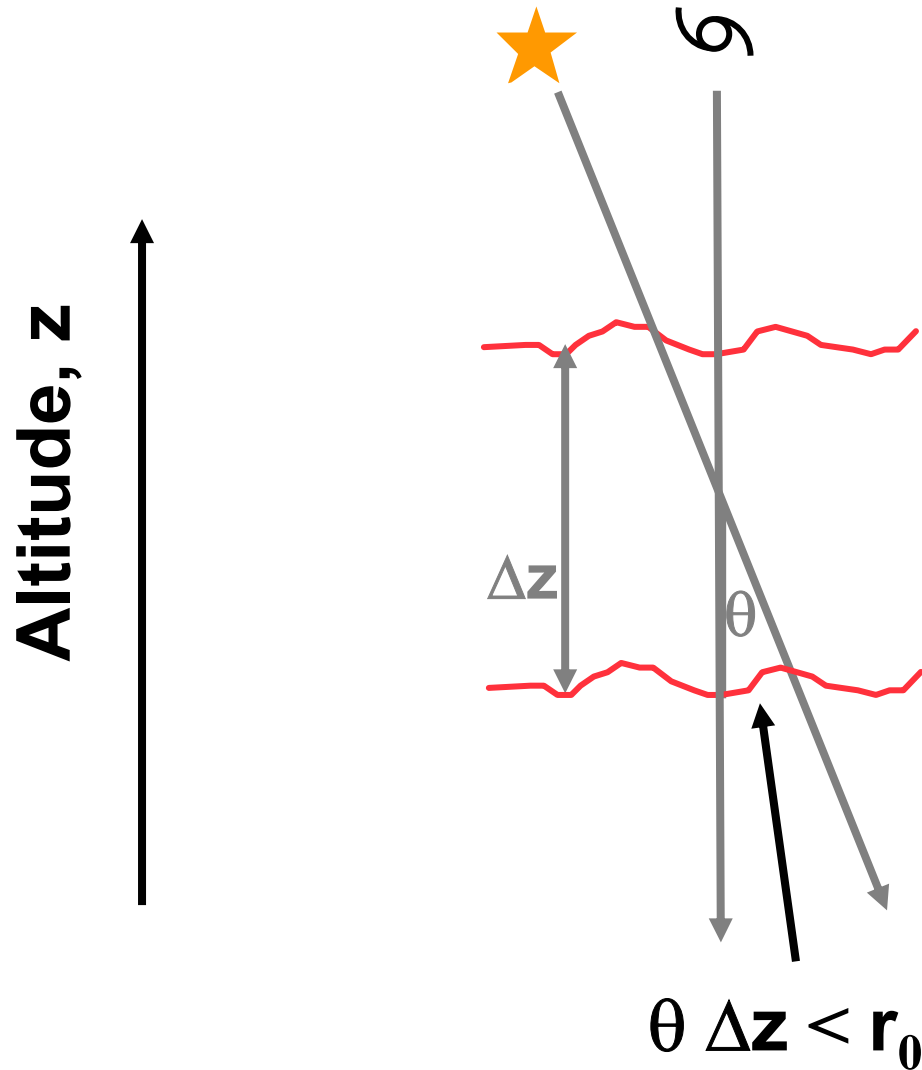
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- **Wider spacing of guide stars gives less dense sampling in k-space.  $\delta z \delta k_z \sim \text{constant}$  so thin  $\delta z$  layers result in wide  $\delta k_z$  lines. Wider  $\delta k_z$  lines (thin layer of turbulence) means fewer samples (guide stars) are required.**
- **Unfortunately, this does not extend to multiple thin layers. E.g. for two thin  $\delta z$  layers spaced apart by  $\Delta z$ , the k-space line is only  $1/\Delta z$  wide, not  $1/\delta z$ . So it doesn't change the situation much from turbulence over the entire  $\Delta z$  region.**
- **Because of the angle of the sample lines in k-space,  $k_z$  sampling interval increases with  $k_x$ , so higher  $k_x$  spatial frequencies are less well measured. The cut off is near  $k_x < 1/[\Delta z(\theta - \theta_{gs})]$**



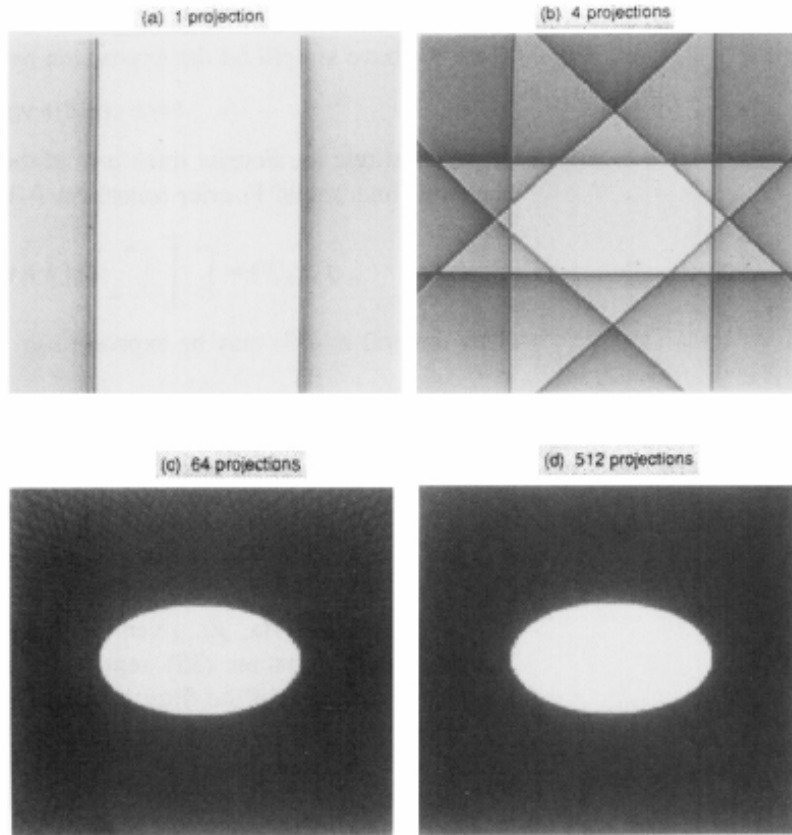
# $k_x < 1/[\Delta z(\theta - \theta_{gs})]$ requirement spatial interpretation





# Back-projection algorithm

Kak, 1988

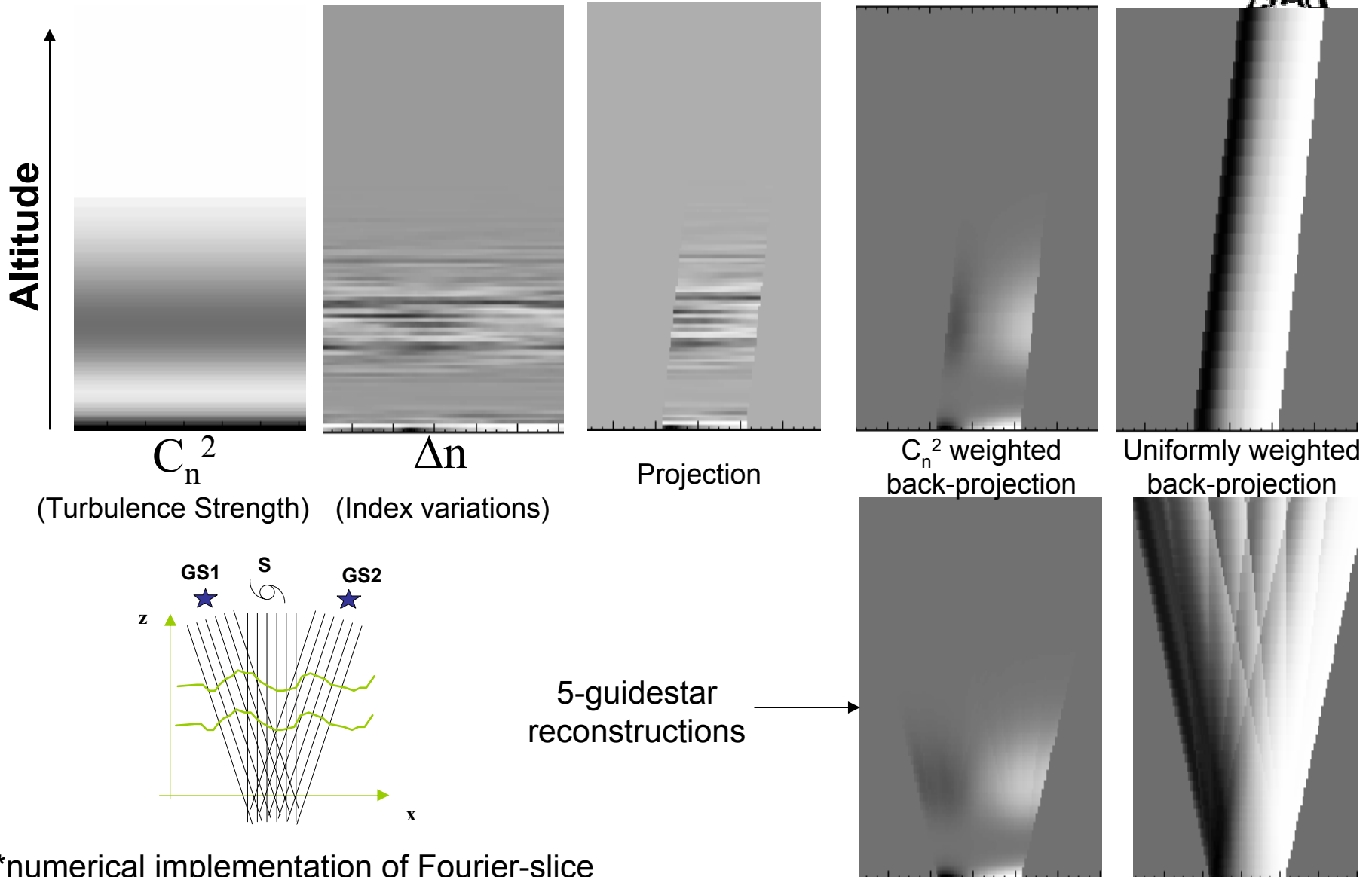


Example tomographic reconstruction of an ellipsoidal target using back projections

- Concepts are borrowed from tomography
- The Fourier-slice reconstruction algorithm is reformulated in terms of projection, filtered back-projection steps (mathematically equivalent to Fourier-slice in the infinite aperture plane wave case).
- *Advantages* are: finite aperture, “fan rays”, finite # of probe directions are brought out explicitly.
- Back propagation is mathematically equivalent to Tokovinin’s algorithm if we take into account  $C_n^2(z)$  profile when back-propagating through the atmosphere. (Proven for  $n_{gs}=1$  case, “statistical addition” must be used for  $n_{gs}>1$ )



# Example NGS finite aperture reconstruction using tomographic back-projection\*

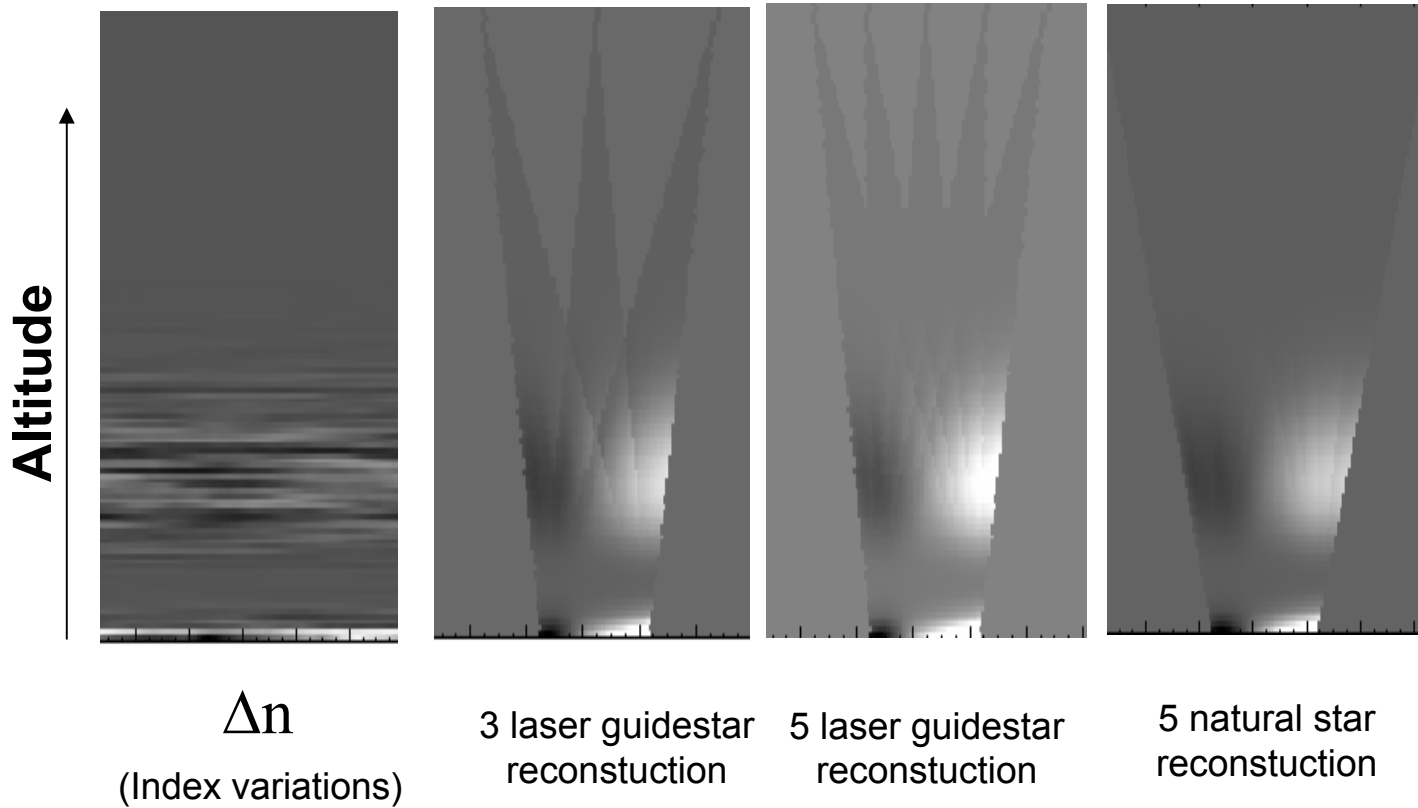
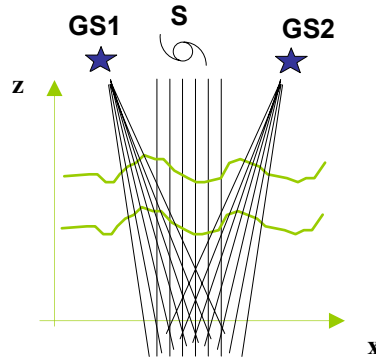


\*numerical implementation of Fourier-slice

note: Tokovinin & Viard's algorithm is equivalent to a  $C_n^2$  weighted back-projection



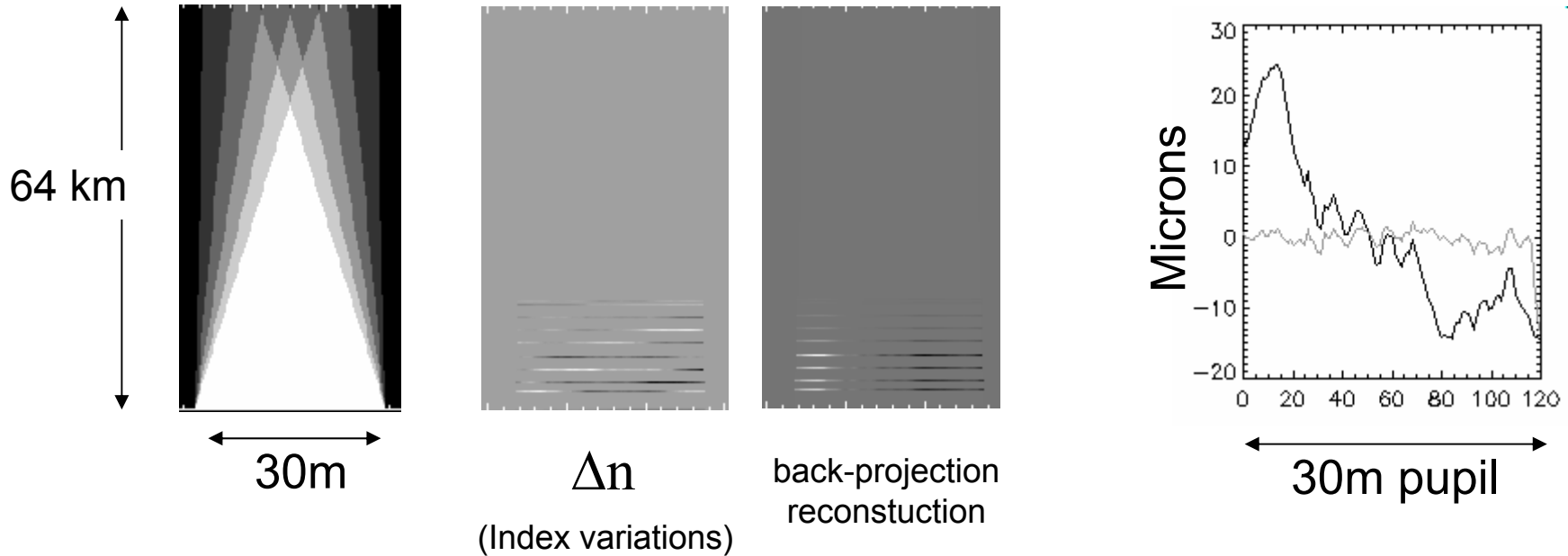
# Cone (laser guide star) back-projections





# LGS tomography error predictions

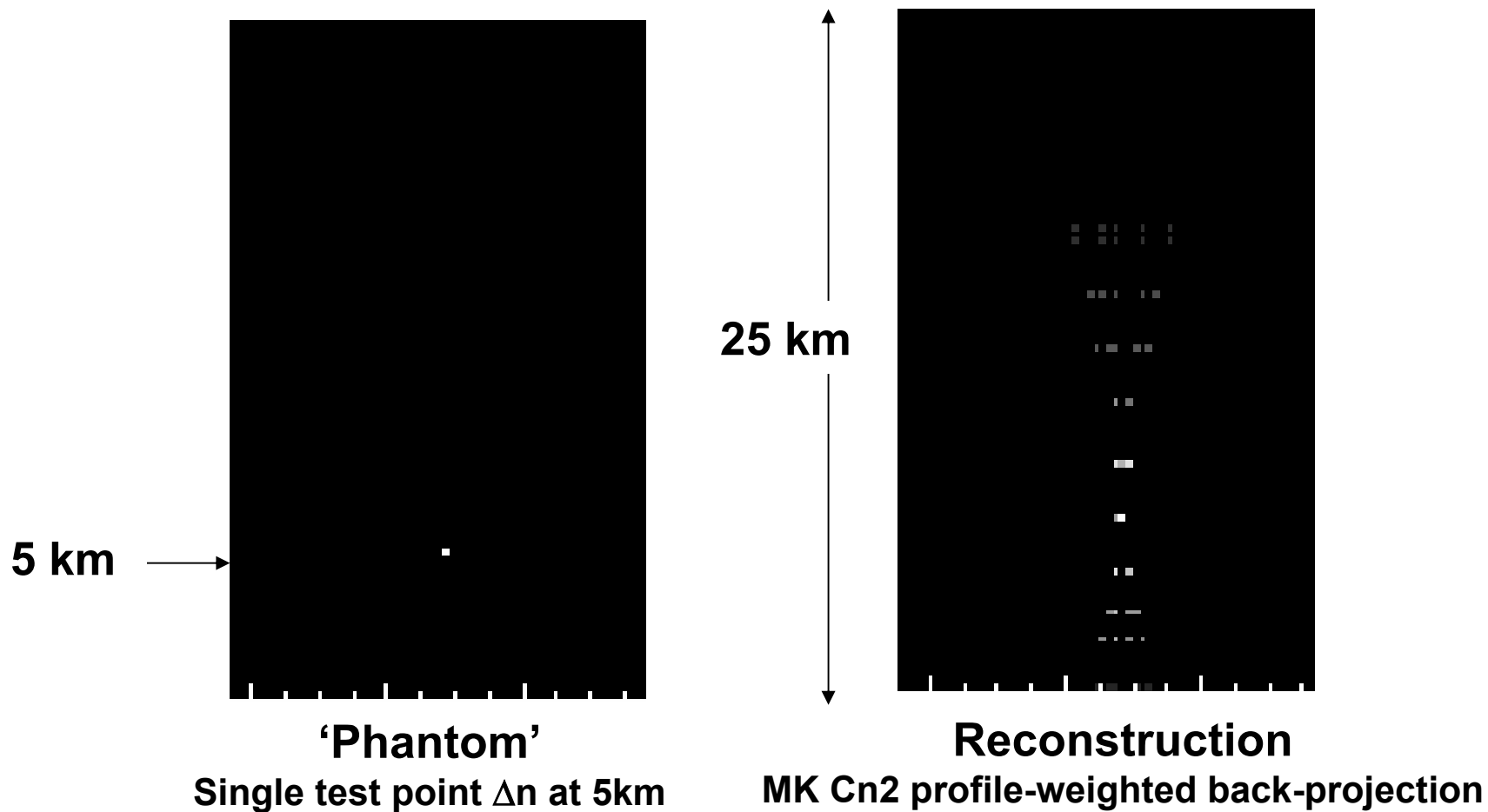
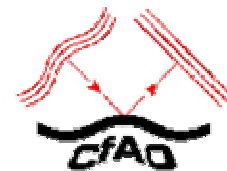
30 m telescope, Mauna Kea 12-layer profile



- Predicts anisoplanatism, cone-effect (validated in single-guidestar simulations)
- Provides insight into number and placement of guidestars
- Back-projection gives ~100x improvement over uncorrected seeing – problem: need ~1000x



# Simulation with a test-point $\Delta n$ demonstrates back-projection's inherent error





# Summary



- Interpretation of MCAO in terms of tomography give insight into system design and performance.
- Includes dependence on number and placement of guidestars.
- Separates the tomography effects from other error sources (e.g. guidestar brightness, DM sampling)
- Predicts “traditional” anisoplanatism and cone effect
- However: back-propagation reconstruction algorithm appears to be 5 to 10 x less accurate than regularized least-squares solutions. Implies back-propagation is not a good real-time reconstructor.